

**Giant Single-Molecule Magnets: A {Mn<sub>84</sub>} Torus and Its Supramolecular Nanotubes\*\***

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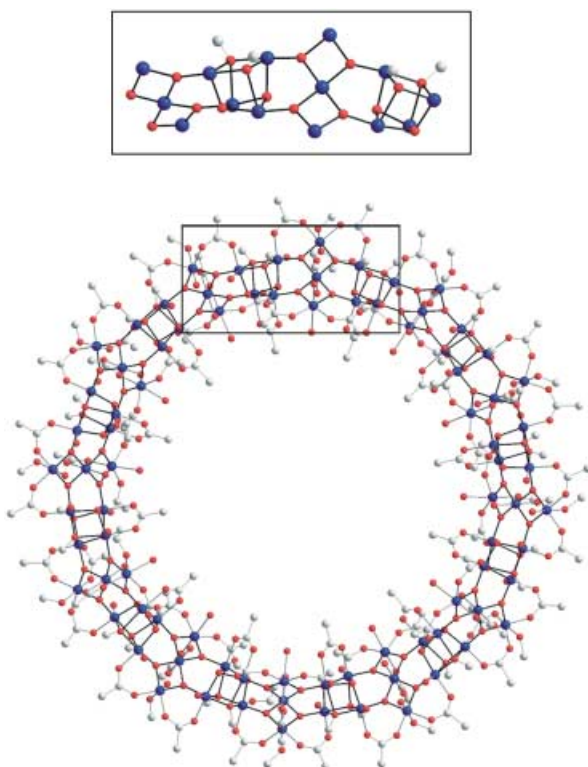
The discovery that individual molecules can function as magnets provided a new, “bottom-up” approach to nanoscale magnetic materials,<sup>[1–3]</sup> and such molecules have since been called single-molecule magnets (SMMs).<sup>[4]</sup> Each molecule is a single-domain magnetic particle that, below its blocking temperature, exhibits the classical macroscale property of a magnet, namely magnetization hysteresis. In addition, SMMs straddle the classical/quantum interface in also displaying quantum tunneling of magnetization (QTM)<sup>[5,6]</sup> and quantum phase interference,<sup>[7]</sup> which are the properties of the micro-scale. SMMs have various potential applications, including very high-density information storage with each bit stored as the magnetization orientation of an individual molecule, and as quantum bits for quantum computing<sup>[8]</sup> by taking advantage of the quantum superposition of states provided by the QTM. For a number of reasons, including facilitating development of techniques for addressing individual SMMs, we have sought to synthesize SMMs of very-large dimensions (by molecular standards). In effect, can the molecular (or bottom-up) approach reach the size regime of the classical (or top-down) approach to nanoscale magnetic materials? Indeed, herein we report a giant Mn<sub>84</sub> SMM. It has a 4 nm diameter torus structure, exhibits both magnetization hysteresis and QTM, and crystallizes as supramolecular nanotubes.

The compound [Mn<sub>84</sub>O<sub>72</sub>(O<sub>2</sub>CMe)<sub>78</sub>(OMe)<sub>24</sub>(MeOH)<sub>12</sub>-(H<sub>2</sub>O)<sub>42</sub>(OH)<sub>6</sub>] $\cdot$ *x*H<sub>2</sub>O $\cdot$ *y*CHCl<sub>3</sub> (**1** $\cdot$ *x*H<sub>2</sub>O $\cdot$ *y*CHCl<sub>3</sub>) was obtained from the reaction of [Mn<sub>12</sub>O<sub>12</sub>(O<sub>2</sub>CMe)<sub>16</sub>(H<sub>2</sub>O)<sub>4</sub>] $\cdot$ 4H<sub>2</sub>O $\cdot$ 2MeCO<sub>2</sub>H (**2**) with (N<sup>n</sup>Bu<sub>4</sub>)(MnO<sub>4</sub>) in MeOH that contained a little acetic acid, followed by filtration and layering of the filtrate with chloroform. After several weeks, the well-formed reddish-brown crystals were isolated in a 20% overall yield, based on the total available Mn equivalents. The compound crystallizes in hexagonal space group *P*6 with the asymmetric unit containing 1/6 of the molecule and approximately 25 water and 2 chloroform solvent mole-

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[\*\*] This work was supported by the U.S. National Science Foundation (Grant CHE-0123603). We thank Philip C. E. Stamp for useful discussions.

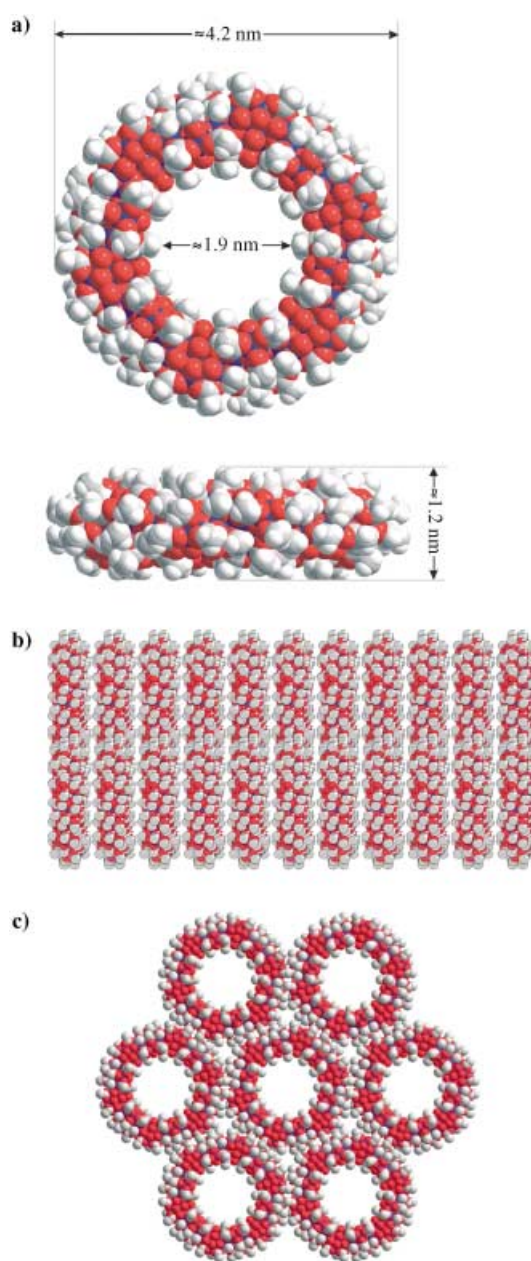
cles of crystallization, which are severely disordered. The structure<sup>[9]</sup> (Figure 1) comprises a  $\{\text{Mn}_{84}\}$  torus with  $C_6$  crystallographic symmetry composed of alternating near-linear  $[\text{Mn}_3\text{O}_4]$  and cubic  $[\text{Mn}_4\text{O}_2(\text{OMe})_2]$  subunits. All the



**Figure 1.** The structure of the  $\{\text{Mn}_{84}\}$  torus, excluding hydrogen atoms. The rectangle shows the repeating  $\{\text{Mn}_{14}\}$  unit that represents the contents of the asymmetric unit; for clarity, it is reproduced above the structure without the carbon atoms (except for those of the  $\text{MeO}^-$  groups to clarify their positions in the  $\{\text{Mn}_{14}\}$  unit). Colour code: Mn blue; O red; C grey.

metal centers are six-coordinate. Close inspection of the Mn–O bond lengths, the Mn bond valence sum calculations (2.82–3.26),<sup>[10]</sup> and the detection of Jahn–Teller axial elongations as expected for a  $d^4$  metal ion in near-octahedral geometry, identified the metal ions as all being in the  $\text{Mn}^{3+}$  oxidation state.

A better appreciation of the structure and the size of the molecule is provided by the space-filling plots of Figure 2, which show that the torus has a diameter of about 4.2 nm and a thickness of about 1.2 nm, with a central hole of diameter 1.9 nm. The  $\text{Mn}_{84}$  molecules order within the crystal in an aesthetically pleasing manner, thus giving nanotubular stacks parallel to the crystal  $c$  axis and to their neighbors (Figures 2b,c). This yields a hexagonal close packing analogous to densely packed straws in a box. The crystalline structure of  $\text{Mn}_{84}$  thus displays extensive cylindrical channel formation along one dimension. Molecules in neighboring chains are exactly adjacent (Figure 2b), and thus the structure may be alternatively described as consisting of graphite-like  $\{\text{Mn}_{84}\}$  sheets lying on top of each other with perfect registry.

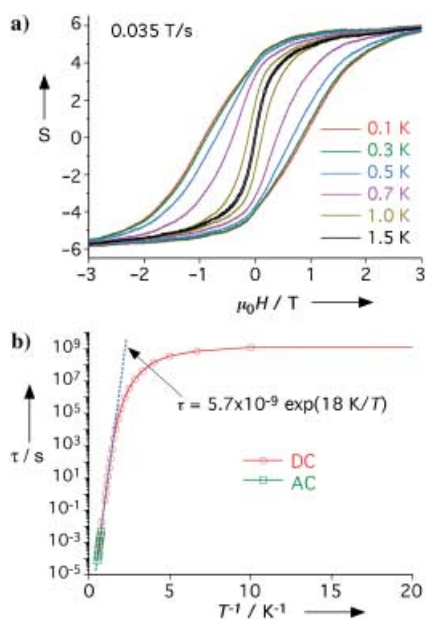


**Figure 2.** Space-filling representations of  $\{\text{Mn}_{84}\}$ , and its supramolecular aggregation into ordered nanotubes and sheets. a) Space-filling representations (including hydrogen atoms) from viewpoints perpendicular (top) and parallel (bottom) to the plane of the torus, showing the dimensions of the molecule and its central hole. b) Ordered arrangement of  $\{\text{Mn}_{84}\}$  molecules (excluding hydrogen atoms) into two adjacent supramolecular nanotubes viewed perpendicular to the propagation axis; this view emphasizes the exact registry of molecules in adjacent tubes and thus the sheetlike structure formed. c) A view along the propagation axes of seven tubes showing the hexagonal packing of neighbors within a single sheet of molecules (excluding hydrogen atoms). Color code: Mn blue; O red; C grey; H white.

The magnetic properties of  $\{\text{Mn}_{84}\}$  have been investigated by both DC and AC methods. The molecule comprises eighty-four  $\text{Mn}^{3+}$  ions ( $S=2$ ), and it was anticipated that there would be fairly strong pairwise exchange interactions between them because they are all monoatomically bridged by  $\text{O}^{2-}$  or  $\text{MeO}^-$

groups. Preliminary magnetic susceptibility ( $\chi_M$ ) studies on polycrystalline samples down to 1.8 K indicated a small uncompensated molecular spin of  $S \approx 6$ . In the presence of sufficient magnetic anisotropy of the easy axis (Ising) type, a spin of this magnitude is sufficient to provide a SMM, and more detailed studies to lower temperatures were therefore carried out to investigate this possibility.

Studies were performed by magnetization measurements on single crystals of  $\{\text{Mn}_{84}\}$  by using an array of micro-SQUIDS (SQUID is superconducting quantum interference device).<sup>[11]</sup> The magnetization versus applied DC field at a  $0.035 \text{ T s}^{-1}$  sweep rate (Figure 3 a) exhibited hysteresis, which



**Figure 3.** The results of magnetic susceptibility studies on single crystals of  $\{\text{Mn}_{84}\}$ . a) Magnetization versus applied DC magnetic field plots, with the field applied along the  $c$  axis (perpendicular to the  $\{\text{Mn}_{84}\}$  torus plane), exhibiting hysteresis loops. The magnetization has been plotted as spin ( $S$ ) per  $\{\text{Mn}_{84}\}$  molecule. A background slope due to low-lying excited states has been subtracted from the data. b) Arrhenius plot constructed by using a combination of out-of-phase AC susceptibility ( $\chi_M''$ ) data and DC magnetization decay data. The dashed line is a fit of the thermally activated region to the Arrhenius relationship; see the text for the fit parameters.

is the diagnostic property of a magnet, thus establishing  $\{\text{Mn}_{84}\}$  as the largest SMM yet discovered. At this sweep rate, the hysteresis becomes evident at 1.5 K, and its coercivity increases with decreasing temperature, as expected for the superparamagnet-like properties of an SMM, before becoming essentially temperature-independent below 0.3 K. The magnetization in Figure 3 a is plotted as spin ( $S$ ) per molecule (determined by using quantitative molar magnetization data). The saturation value indicates a molecular ground state spin of  $S = 6$ , thus confirming the preliminary estimate above. The hysteresis loop shows no sign of the steps diagnostic of QTM that are visible in the hysteresis loops of several smaller SMMs such as  $\{\text{Mn}_{12}\}$ ,<sup>[5,6]</sup>  $\{\text{Mn}_4\}$ ,<sup>[12,13]</sup> and  $\{\text{Fe}_8\}$ ,<sup>[7,14]</sup> and the exchange-biased  $(\{\text{Mn}_4\})_2$  dimer.<sup>[15]</sup> In previous cases in which no steps are visible, for large SMMs such as the  $\{\text{Mn}_{18}\}$ <sup>[2+][16]</sup>

and  $\{\text{Mn}_{30}\}$ <sup>[17]</sup> SMMs, this lack of visible steps is due to a broadening and smearing of the steps from low-lying excited states and a distribution of molecular environments (and thus a distribution of relaxation barriers) caused by disordered lattice solvent molecules and ligand disorder; the magnetic properties of SMMs are sensitive to such relatively small variations in local environments. For  $\{\text{Mn}_{84}\}$ , the large numbers of disordered solvent molecules in the central cavity readily rationalize a distribution of molecular environments.

With Figure 3 a establishing  $\{\text{Mn}_{84}\}$  as an SMM but not showing steps, it became of importance to establish by an alternative means whether quantum tunneling is still occurring in such a giant SMM. This was addressed by collecting magnetization-relaxation data from 1) out-of-phase AC-susceptibility ( $\chi_M''$ ) data at different AC frequencies, and 2) plots of magnetization versus time decay. Both methods provide data of the rate of magnetization relaxation ( $1/\tau$ ) versus  $T$  ( $\tau$  is the relaxation time). These data are plotted as  $\tau$  versus  $1/T$  in Figure 3 b. There is a thermally activated region above about 0.5 K, and fitting of this region to the Arrhenius relationship of Equation (1) gives  $U_{\text{eff}} = 18 \text{ K}$  and  $\tau_0 = 5.7 \times 10^{-9} \text{ s}$ , in which

$$(1/\tau) = (1/\tau_0) \exp(-U_{\text{eff}}/kT) \quad (1)$$

$U_{\text{eff}}$  is the effective (mean) relaxation barrier and  $\tau_0$  is the pre-exponential factor.

The relaxation rate levels off below 0.5 K, and below about 0.2 K it is temperature-independent, which is consistent with relaxation only by ground-state tunneling between the lowest energy  $M_S = \pm 6$  levels of the  $S = 6$  manifold. Such temperature-independent relaxation rates arising from QTM have been observed previously for several other SMMs.<sup>[12–14,16–18]</sup> Independent confirmation of QTM in  $\{\text{Mn}_{84}\}$  was obtained from “quantum hole digging”,<sup>[19]</sup> which can establish resonant QTM even when no steps are apparent in the hysteresis loops. It is thus clear that  $\{\text{Mn}_{84}\}$  still exhibits the quantum behavior that has become a common feature of the much smaller SMMs.

The above results establish  $\{\text{Mn}_{84}\}$  as a giant SMM. In general, SMMs have many advantages over classical magnetic particles made of Co metal,  $\text{Fe}_3\text{O}_4$ , etc., including monodispersity, crystallization as highly ordered ensembles, true solubility (rather than colloidal suspensions), and a shell of organic groups that protects the magnetic cores from surface variations and prevents their contact with each other. The discovery of a  $\{\text{Mn}_{84}\}$  SMM establishes that these advantages can also be extended to this giant molecule, which is still soluble, stable, and crystalline. Such large wheel-like molecules are with precedent, but only in molybdenum chemistry,<sup>[20]</sup> with the largest currently known being the  $\{\text{Mo}_{154}\}$ <sup>[21]</sup> and  $\{\text{Mo}_{176}\}$ <sup>[22]</sup> compounds prepared by Müller and co-workers. Although they contain more metal atoms than  $\{\text{Mn}_{84}\}$ , the  $\{\text{Mo}_{154}\}$  and  $\{\text{Mo}_{176}\}$  wheels have diameters of about 3.4 and about 4.1 nm, respectively. These, and other giant polyoxometallates,<sup>[23–25]</sup> are diamagnetic or nearly so, and none of them exhibit SMM properties.

But although  $\{\text{Mn}_{84}\}$  is very large by molecular standards, it is pertinent to ask how it compares to classical nano-

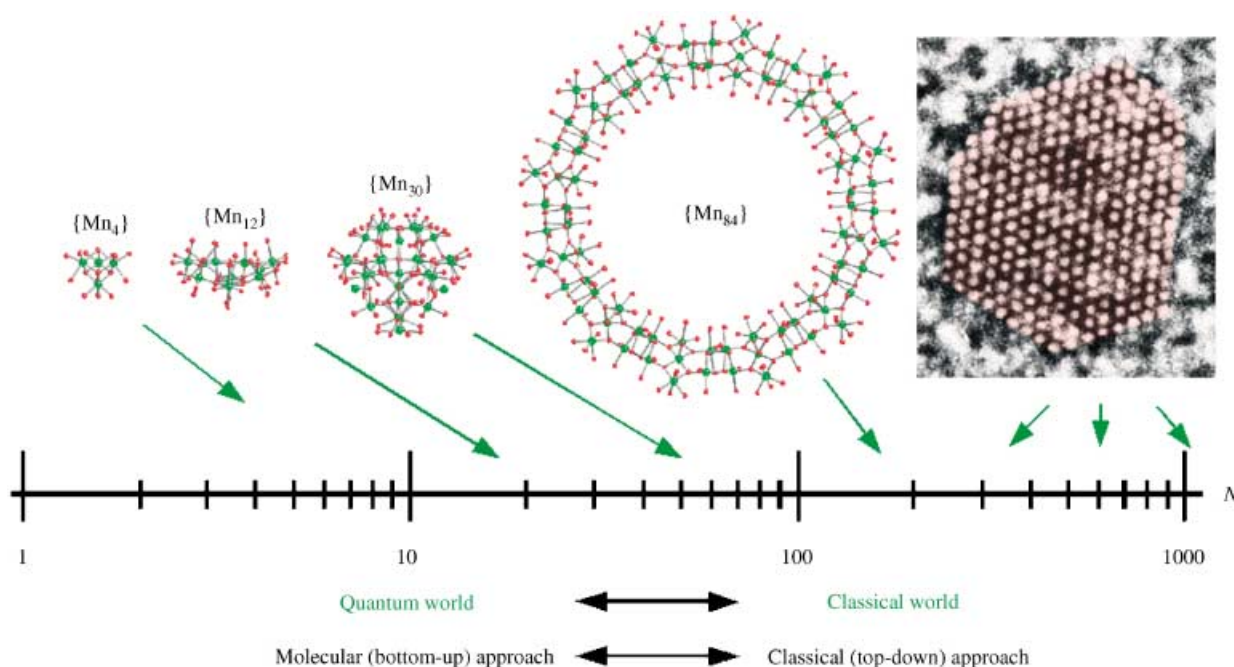
particles. In Figure 4 the sizes of  $\{\text{Mn}_4\}$ ,  $\{\text{Mn}_{12}\}$ ,  $\{\text{Mn}_{30}\}$ , and  $\{\text{Mn}_{84}\}$  SMMs are compared with that of a 3 nm Co nanoparticle recently reported,<sup>[26]</sup> all drawn to the same scale. With a 4.2 nm diameter,  $\{\text{Mn}_{84}\}$  is thus of comparable size to the smallest nanoparticles. The total number of atoms in  $\{\text{Mn}_{84}\}$  is 1032, roughly the same as the 3 nm Co nanoparticle, which contains about 1000 Co atoms. Of course, the former has a very different shape, given its central hole and essentially wheel-like rather than spherical structure. Another useful way, particularly in the physics literature, of comparing the “size” of magnetic systems is by their Néel vector ( $N$ , the sum of the individual spins), which are 7.5, 22, 61, and 168 for  $\{\text{Mn}_4\}$ ,  $\{\text{Mn}_{12}\}$ ,  $\{\text{Mn}_{30}\}$ , and  $\{\text{Mn}_{84}\}$ , respectively. This is the scale used in Figure 4. With a value of 168,  $\{\text{Mn}_{84}\}$  is a far larger spin system than any other molecular cluster (greater than the next largest, 75, observed previously for an  $\text{Fe}_{30}$  species with  $S=0$ <sup>[24,25]</sup>) and is at the lower limit of values found for classical nanoparticles, which can range from a few hundred to many thousands, depending on the precise size and constituent metal; for the 3 nm Co nanoparticle shown, the Néel vector is approximately 1000.

The new  $\{\text{Mn}_{84}\}$  SMM thus essentially represents a long-sought-after meeting of the bottom-up and top-down approaches to nanoscale magnetic materials. It will be interesting to see whether even larger SMMs that contain more metal atoms can be obtained in the near future; we see no reason why this should not be so. It is also interesting that  $\{\text{Mn}_{84}\}$  exhibits magnetization quantum tunneling, a point of relevance to the general question of how large magnetic

nanoparticles can become and still display quantum effects.<sup>[27]</sup> Such studies for classical magnetic nanoparticles have been hampered by complications from distributions of particle size, shape, surface roughness, and spin<sup>[11,28]</sup> a problem that does not exist for crystalline, monodisperse SMMs. Finally, we recognize that the supramolecular architecture of the  $\{\text{Mn}_{84}\}$  molecules as ordered nanotubular stacks offers a variety of future possibilities for accessing interesting new materials, including the insertion of guest molecules (either organic or inorganic, diamagnetic or paramagnetic) or chains (either conducting or insulating) into or through the tubes. In addition, the large shape anisotropy and solubility of  $\{\text{Mn}_{84}\}$  suggest that it should be fairly straightforward to deposit and visualize these molecular rings on surfaces.

### Experimental Section

**Compound 1:** A freshly-prepared solution of  $\text{NnBu}_4\text{MnO}_4$  (0.30 g, 0.83 mmol) in MeOH (10 mL) and glacial acetic acid (0.75 ml) was added to a slurry of  $[\text{Mn}_{12}\text{O}_{12}(\text{O}_2\text{CMe})_{16}(\text{H}_2\text{O})_4]\cdot 4\text{H}_2\text{O}\cdot 2\text{HO}_2\text{CMe}$  (0.425 g, 0.206 mmol) in MeOH (15 mL) over a period of two minutes. The mixture was left under magnetic stirring for a few minutes, and then filtered to give a reddish-brown filtrate. Equal amounts of the filtrate were placed in five different vials and layered with chloroform. After a few weeks reddish-brown crystals appeared, and were left to grow for several more weeks to give well shaped X-rays quality crystals. When product formation was deemed to be complete, the crystals were collected by filtration, washed with chloroform and dried in vacuum. The yield was 0.10 g (20% based on total Mn). Elemental analysis (%) calcd for  $\text{C}_{192}\text{H}_{464}\text{O}_{322}\text{Mn}_{84}$



**Figure 4.** The position of  $\{\text{Mn}_{84}\}$  on a size scale spanning atomic to nanoscale dimensions. On the far right is shown a high-resolution transmission electron microscopy view along a  $[110]$  direction of a typical 3 nm diameter cobalt nanoparticle exhibiting a face-centered cubic structure and containing about 1000 Co atoms.<sup>[26]</sup> The  $\{\text{Mn}_{84}\}$  molecule is a 4.2 nm diameter particle. Also shown for comparison are the indicated smaller Mn SMMs, which are drawn to scale. An alternative means of comparison is the Néel vector ( $N$ ), which is the scale shown. The green arrows indicate the magnitude of the Néel vectors for the indicated SMMs, which are 7.5, 22, 61, and 168 for  $\{\text{Mn}_4\}$ ,  $\{\text{Mn}_{12}\}$ ,  $\{\text{Mn}_{30}\}$  and  $\{\text{Mn}_{84}\}$ , respectively. The green arrows from the Co nanoparticle are merely meant to indicate that the Néel vector of nanoparticles can take many values, depending on the exact size and the identity of the constituent metal.



(1·10H<sub>2</sub>O): C 18.39, H 3.73; found: C 18.35, H 3.45. Selected IR data (KBr pellet):  $\tilde{\nu}$  = 3400(s, br), 2965(w), 2931(w), 1576(s, br), 1527(s, br), 1425(s, br), 1348(m), 1250(w), 1028(m), 949(w), 696(m), 666(s), 623(s), 580(s), 548(s), 506(m).

Received: November 18, 2003 [Z53352]

Published Online: March 9, 2004

**Keywords:** cluster compounds · magnetic properties · manganese · nanomaterials · single-molecule magnets

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- [9] a) Crystal structure data for **1**: C<sub>192</sub>H<sub>444</sub>O<sub>312</sub>Mn<sub>84</sub>,  $M_r = 12360.24$ , hexagonal, P6,  $a = 41.345(2)$ ,  $c = 14.180(2)$  Å,  $V = 20992(3)$  Å<sup>3</sup>,  $T = 123$  K,  $Z = 1$ ,  $d_{\text{calc}} = 1.305$  g cm<sup>-3</sup>, 85 636 reflections collected, 16 019 unique ( $R_{\text{av}} = 0.0767$ ),  $R1 = 0.0849$ ,  $wR2 = 0.2117$ , using 12 520 reflections with  $I > 2\sigma(I)$ . The asymmetric unit, in addition to 1/6th of the {Mn<sub>84</sub>} cluster, also contains an estimated 25 water and two chloroform molecules of crystallization. Most of the solvent molecules are disordered and could not be modeled properly, thus the program SQUEEZE, a part of the PLATON package<sup>[9b]</sup> of crystallographic software, was used to calculate the solvent disorder area and remove its contribution to the overall intensity data. CCDC-224300 contains the supplementary crystallographic data for this paper. These data can be obtained free of charge via [www.ccdc.cam.ac.uk/conts/retrieving.html](http://www.ccdc.cam.ac.uk/conts/retrieving.html) (or from the Cambridge Crystallographic Data Centre, 12 Union Road, Cambridge CB21EZ, UK; fax: (+44) 1223-336-033; or [deposit@ccdc.cam.ac.uk](mailto:deposit@ccdc.cam.ac.uk)); b) PLATON: A. L. Spek, *Acta Crystallogr. Sect. A* **1990**, 46, 1–34.
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